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# Ka-Band Linear to Circular Polarization Converter Based on Multilayer Slab With Broadband Performance

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**ABSTRACT** In this paper, a Ka-band polarization converter is presented, which is based on multilayer slab. In order to improve impedance matching, metallic circular traces are printed periodically on each dielectric multilayer slab. Simulated results of the polarizer show that it can transform linearly polarized (LP) to circularly polarized (CP) fields over a frequency band from 23 to 35GHz (42%) with an insertion loss less than 0.5 dB. The transmitted CP wave by the polarizer is approximately robust under oblique illuminations. The polarizer is fabricated and measured by a wideband horn antenna satisfying the simulated results. Next, in order to design a high-gain CP structure around 30 GHz, an 8-element LP array antenna with Chebyshev tapered distribution is designed and integrated with the polarizer. Obviously, the antenna limits the overall bandwidth (nearly 28 to 31.5 GHz) due to the narrowband nature of the LP antenna array. When the polarizer is illuminated by an incident LP wave, the two linear components of the transmitted wave with approximately equal amplitudes and 90° phase difference on the frequency band of interest are produced. Experimental results of the proposed structure show a pure CP with a gain of 13 dBi at 30 GHz, which can be suitable for millimeter wave communication.

**INDEX TERMS** Polarizer, Ka-band, linear, circular, polarization converter, dielectric multilayer.

## I. INTRODUCTION

Polarization properties of antennas play a critical role in wireless communication systems. Recently, because of the advanced signal propagation properties, systems using CP antennas offer significant performance benefits over systems using traditional linearly polarized antennas. The use of CP is considered as a solution for addressing challenges associated with adverse weather conditions, and polarization mismatch, none-line-of-sight applications, phasing issues, and multipath effects [1]–[3]. Various CP antennas have been presented in the literature in the last few decades. A suitable way of creating a CP wave is to radiate a linearly polarized electromagnetic wave to a polarizer plane. This method has become appealing in situations where the radiating structure contains a planar array. For this case, the polarizer is placed on top of the antenna to transmit CP waves. Basically, the polarization converter decomposes the LP incident plane wave into two orthogonal components

of nearly equal amplitude with a phase difference of 90° to generate a CP transmitted plane wave. A transmission-type polarization converters, which are also named polarizers, have been extensively investigated by researchers in the open literature for several applications [4]–[13]. For example, some polarizers were designed utilizing frequency-selective surface (FSS) structures [4]–[7], and others are based on traditional hybrid meander line and loop configuration [8]–[10]. In [4] a multilayer FSS was designed using square patches with truncated corners to obtain 90° phase difference between orthogonal linear components. The polarizer presented in [5] employs a miniaturized-element FSS composed of arrays of subwavelength capacitive patches and inductive wire grids separated by dielectric substrates. The polarizer in [6] uses a double-sided partially reflective surface (PRS), which controls independently the reflection and transmission coefficients. The work in [7] suggested a CP converter utilizing a four layer FSS consists of split rings

bisected by a metal strip, in which distance between different layers is a quarter wavelength, but the layers' dimensions are different. Although structures suggested in [4]–[7] are broad band, they consist of multilayer structures with different configurations and dimensions for each layer, ending up with complicated geometry and configuration for the polarizers. In [9] single layer polarizer based on hybrid meander line and loop geometry are proposed. This structure can cover the AR bandwidth of 46.8%. However, the peak insertion loss is 3dB. In order to achieve a CP in [10], a meander line polarizer has been employed with the linear SIW antenna, while the structure does not have broadband behavior. Another approach used in [11] included rotating FSS structure that was implemented using the substrate integrated waveguide (SIW) technology. The rotating FSS in [12] utilizing a triple-mode SIW cavity resonator resulted in improving the transmission bandwidth. Meta-surface and metamaterial have been applied in [12]–[13]. A meta-surface of 16 rectangular loop unit cells each with a diagonal microstrip line is positioned above a source antenna in [12]. An ultra-compact chiral metamaterial (CMM) using triple-layer twisted split-ring resonators (TSRRs) structure was suggested in [13]. Despite the benefits proposed by these designs, the majority of them employ resonant elements that tend to be narrowband. In this study, a Ka-band transmission-type polarizer based on multilayer dielectric structure to convert LP incident electromagnetic wave into a CP transmitted wave is proposed. To simplify the structure, a unit cell is used in all dielectric layers. In order to illuminate the polarizer, an 8-element linearly polarized linear array antenna with Chebyshev tapered distribution to control the radiated pattern side lobe level (SLL) is also designed. In the configuration, the antenna is placed at a half free space wavelength from the polarizer with a 45° rotation. The size of the polarizer is determined using a diffraction and transmission rays' model. The benefits of the proposed structure include (a) wideband polarizer with high CP purity, (b) polarizer can function for broad-angle oblique incident waves, (c) less insertion loss, (d) simple structure using same unit cell in all layers, (e) employing a low SLL linear array antenna, and (f) implementation at mm-wave Ka-band. The proposed polarizer is based on a nano-polarizer with periodic and two-dimensional dielectric high-contrast gratings [14]. The basic theory and operation of the polarizer are discussed in Section II. Section III outlines the design process of the polarizer's unit cell. Simulated results of the proposed polarizer are presented and discussed in Section IV. The design of single and 8-element linear array antenna with Chebyshev tapered distribution are described in Section V. Section VI presents the measured results of the combined structure of the antenna and polarizer. Finally, a brief conclusion is given in Section VII.

## II. THEORETICAL ANALYSIS

In this section, we consider the theoretical analysis of the polarizer. Fig. 1 exhibits the topology of the polarizer which converts linearly polarized EM wave into a circular one.

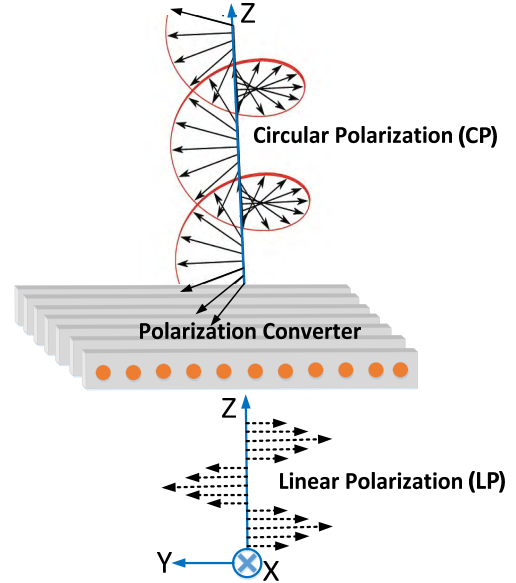


FIGURE 1. Schematic model of a linear-to-circular polarizer.

Let us consider an incoming plane wave that propagates in the +z direction, whose incident and transmitted electric fields can be expressed as [15]

$$E^{inc}(x, y, z, t) = \begin{pmatrix} E_x^{inc} \\ E_y^{inc} \end{pmatrix} e^{i(kz - t)} \quad (1)$$

Where  $\omega$ ,  $k$ ,  $E_x$ , and  $E_y$  are the frequency, wave number, and the complex amplitudes of x- and y- electric field components, respectively [29].

$$E^{tr}(x, y, z, t) = \begin{pmatrix} E_x^{tr} \\ E_y^{tr} \end{pmatrix} e^{i(kz - t)} \quad (2)$$

To better understand the polarization conversion of the polarizer, the transmission coefficient T.C. is defined as:

$$T.C. = \frac{E^{tr}}{E^{inc}} \quad (3)$$

The linear transmission matrix can be expressed as [29]

$$\begin{pmatrix} E_x^{tr} \\ E_y^{tr} \end{pmatrix} = (T.C.)_{Linear} \begin{pmatrix} E_x^{inc} \\ E_y^{inc} \end{pmatrix} = \begin{pmatrix} T_{XX} & T_{XY} \\ T_{YX} & T_{YY} \end{pmatrix} \begin{pmatrix} E_x^{inc} \\ E_y^{inc} \end{pmatrix} \quad (4)$$

Thus, the circularly polarized transmitted EM wave can be defined as [29]:

$$\begin{pmatrix} E_{RHCP}^{tr} \\ E_{LHCP}^{tr} \end{pmatrix} = \begin{pmatrix} E_x^{tr} + jE_y^{tr} \\ E_x^{tr} - jE_y^{tr} \end{pmatrix} = (T.C.)_{Circular} \begin{pmatrix} E_x^{inc} \\ E_y^{inc} \end{pmatrix} \quad (5)$$

$$\begin{aligned} (T.C.)_{Circular} &= \begin{pmatrix} T_{RHCP\_X} & T_{RHCP\_Y} \\ T_{LHCP\_X} & T_{LHCP\_Y} \end{pmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{pmatrix} T_{XX} + jT_{YX} & T_{XY} + jT_{YY} \\ T_{XX} - jT_{YX} & T_{XY} - jT_{YY} \end{pmatrix} \end{aligned} \quad (6)$$

where the factor  $\frac{1}{\sqrt{2}}$  in (6) is due to power normalization. The matrix in (6) states the ability of the polarizer to convert an LP electromagnetic wave into a CP one. For the proposed structure, the cross transmission coefficients are considered to be negligible ( $T_{YX} \approx T_{XY} \approx 0$ ). Therefore, we have

$$E_{RHCP}^t = \frac{1}{\sqrt{2}} \left[ (T_{XX}) E_x^{inc} + (jT_{YY}) E_y^{inc} \right] \quad (7)$$

$$E_{LHCP}^t = \frac{1}{\sqrt{2}} \left[ (T_{XX}) E_x^{inc} - (jT_{YY}) E_y^{inc} \right] \quad (8)$$

In order to generate a CP wave within a desirable frequency band, the following requirements should be satisfied.

$$|T_{xx}| = |T_{yy}| \quad (9)$$

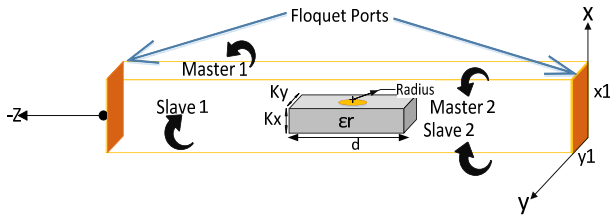
$$\Delta\theta = \theta_{xx} - \theta_{yy} = \pm \frac{\pi}{2} \quad (10)$$

In (10),  $+\frac{\pi}{2}$  corresponds to a RHCP transmitted wave, while  $-\frac{\pi}{2}$  represents a LHCP one. The axial ratio (AR) of the polarizer is evaluated utilizing transmission line theory [16]:

$$AR = \left( \frac{|T_{xx}|^2 + |T_{yy}|^2 + \sqrt{a}}{|T_{xx}|^2 + |T_{yy}|^2 - \sqrt{a}} \right)^{\frac{1}{2}} \quad (11)$$

Where

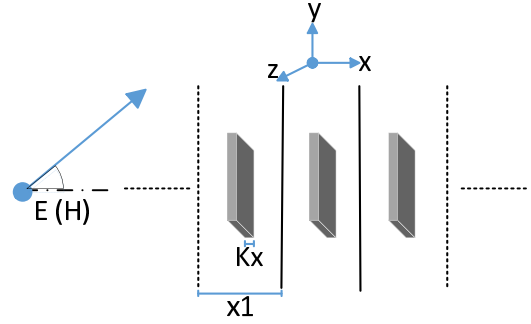
$$a = |T_{xx}|^4 + |T_{yy}|^4 + 2|T_{xx}|^2|T_{yy}|^2 \cos(2\Delta\theta) \quad (12)$$



**FIGURE 2.** The unit cell model by HFSS software ( $x_1 = 1.6\text{mm}$ ,  $y_1 = 1\text{mm}$ ,  $d=4.6\text{mm}$ ,  $k_x=0.254\text{mm}$ ,  $k_y=1\text{mm}$ ,  $\epsilon_r=10.2$ ,  $\text{Radius}=0.3\text{mm}$ ).

### III. POLARIZER DESIGN

The schematic model of the proposed multilayer dielectric LP to CP polarizer is shown in Fig. 1. The polarizer is illuminated by a planar LP electromagnetic wave propagating in the  $k$  direction and tilted away from the  $x$ -axis by  $45^\circ$ . Therefore, the electric field is decomposed into two orthogonal components: the former along the  $x$ -axis and the latter along the  $y$ -axis. Fig. 2 exhibits the unit cell model simulated by ANSYS HFSS<sup>TM</sup>. The unit cell consists of two Floquet ports in distance of one wavelength ( $\lambda$ ) far from the dielectric with constant  $\epsilon_r = 10.2$  and thickness  $k_x = 0.254\text{mm}$ . The width and length of the dielectric piece are  $k_y=1\text{mm}$  and  $d=4.6\text{mm}$ , respectively. The radius of the metal circle is  $0.3\text{mm}$ . In addition, the method applied in the unit cell model is the periodic boundary conditions of the Master/Slave in order to evaluate the infinite structure. Both Floquet ports contain two dominant modes with two orthogonal waves along  $x$ - and  $y$ -axis. The main idea of the LP-to-CP



**FIGURE 3.** One dimensional lattice of dielectric slabs of width  $k_x$  in a periodic lattice with period  $x_1$ .

polarizer and simulation model have come from [14], [17], and [18]. In order to study a fundamental theory for multilayer dielectric slabs, Fig. 3 is presented. A representative one-dimensional periodic array of dielectric slabs with period  $x_1$  and dielectric insert width  $K_x$  is depicted in Fig. 3. If the electric field does not have a component in the  $x$ -direction, the mode is denoted  $\text{TE}_x$  or horizontal. If the magnetic field does not have a component in the  $x$ -direction, the mode is denoted  $\text{TM}_x$  or vertical.

The effective dielectric constants given in two directions by the volumetric average of the constitutive phases are [19], [20].

$$\epsilon_{rx} = \begin{cases} \epsilon_r, & -\frac{k_x}{2} < x < \frac{k_x}{2} \\ 1, & \text{otherwise;} \end{cases} \quad (13)$$

$$\epsilon_{ry} = [1 + (\epsilon_r - 1) \frac{k_x}{x_1}] \quad (14)$$

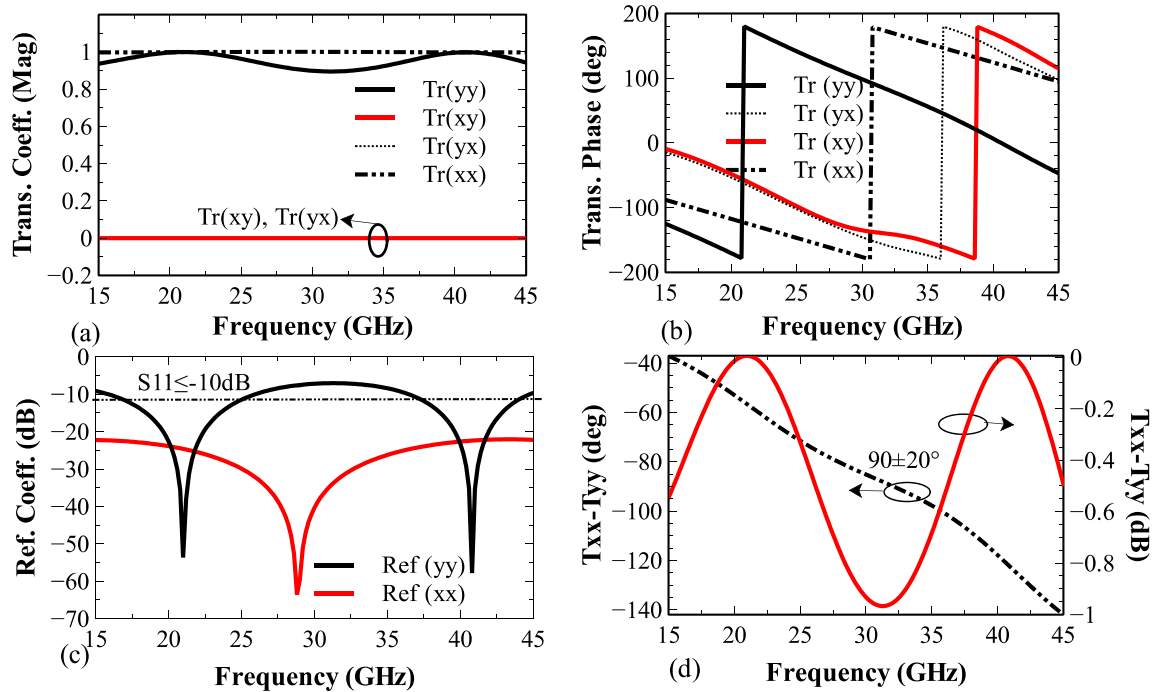
From (13) and (14) it can be noted that the effective dielectric constant in  $y$ -direction depends on  $\frac{k_x}{x_1}$ , while the  $x$ -directed component has a constant value close to 1. Therefore, due to the existence of the periodic dielectric slabs, a phase difference happens between the  $E$ -field orthogonal components when the incident EM wave transmits through the LP-to-CP polarizer. The phase difference is estimated as:

$$\text{Phase Difference} = (\beta_y - \beta_x)d = \frac{2\pi}{\lambda} (\sqrt{\epsilon_{ry}} - \sqrt{\epsilon_{rx}})d \quad (15)$$

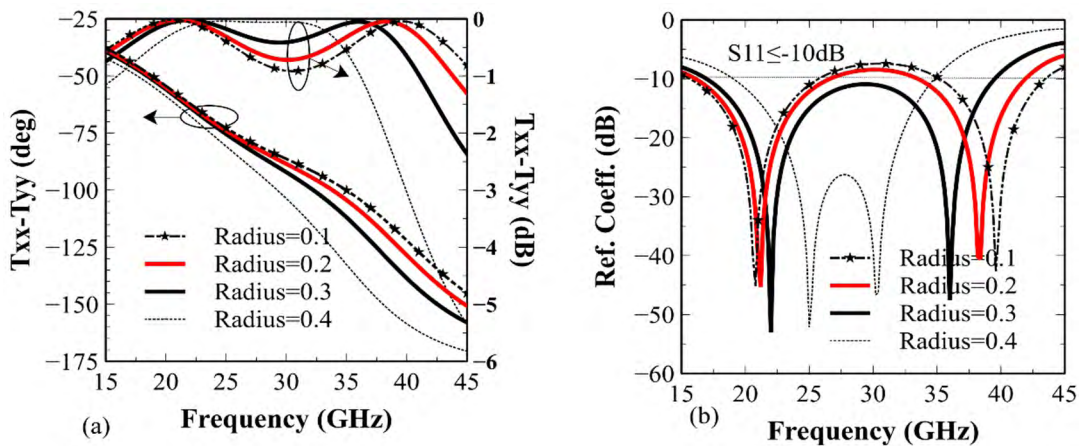
where  $\lambda$ ,  $\beta$  are the wavelength and phase constant, respectively. The phase difference between the  $E$ -field orthogonal components is tuned using a length parameter  $d$ , as shown in Fig. 2. In order to satisfy the CP requirement in (15), the phase difference must be  $\frac{\pi}{2}$ , therefore

$$d = \frac{\lambda_0}{4 (\sqrt{\epsilon_{ry}} - \sqrt{\epsilon_{rx}})} \quad (16)$$

The dielectric multilayer slab characteristics at a resonant frequency ( $f_r$ ) are governed by (13) through (16), in which  $\epsilon_{rx}$  can be estimated to 1. Thus, at 30 GHz with dielectric constant 10.2, thickness  $k_x=0.254\text{mm}$ , and the periodic distance between slabs  $x_1 = 1.6\text{mm}$ ,  $d=4.6\text{mm}$  is obtained.



**FIGURE 4.** The simulated of transmission coefficients of dielectric slab without the conductor in unit cell model (a) magnitude of transmission coefficients, (b) phase of transmission coefficients, (c) reflection coefficient, and (d) difference of phase and magnitude between transmission coefficients.



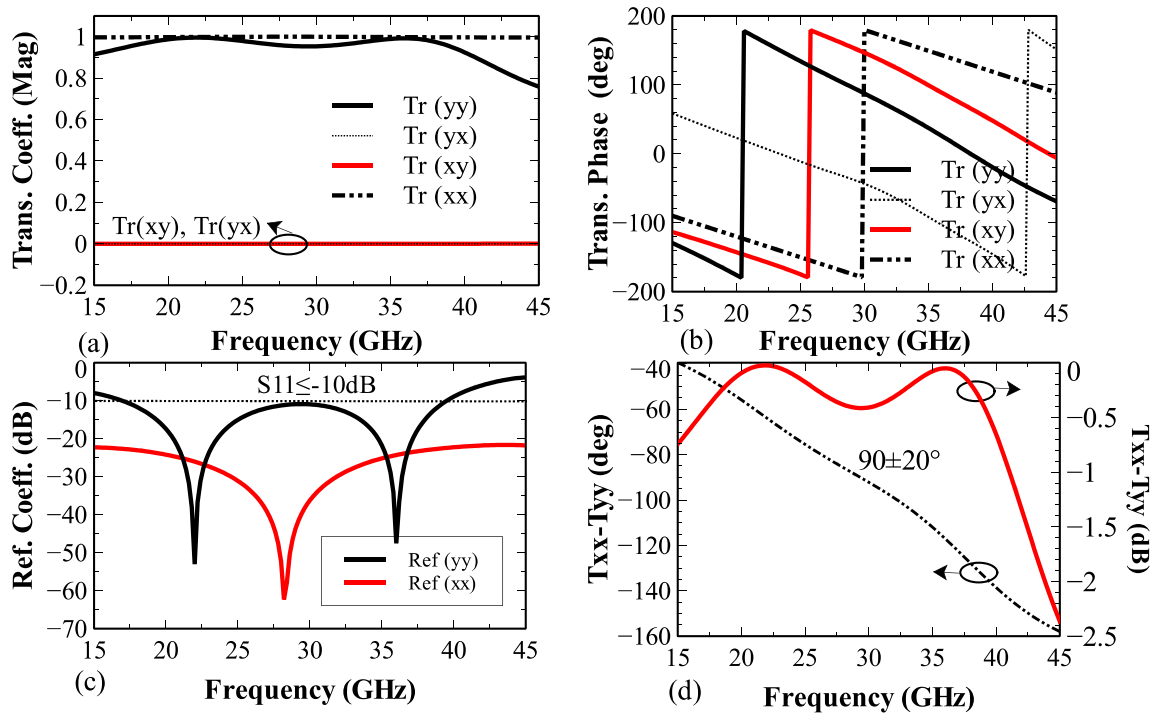
**FIGURE 5.** (a) variations of transmission coefficients and (b) reflection coefficients in the unit cell model when radius of the circle conductor changes.

#### IV. SIMULATION RESULTS OF POLARIZER

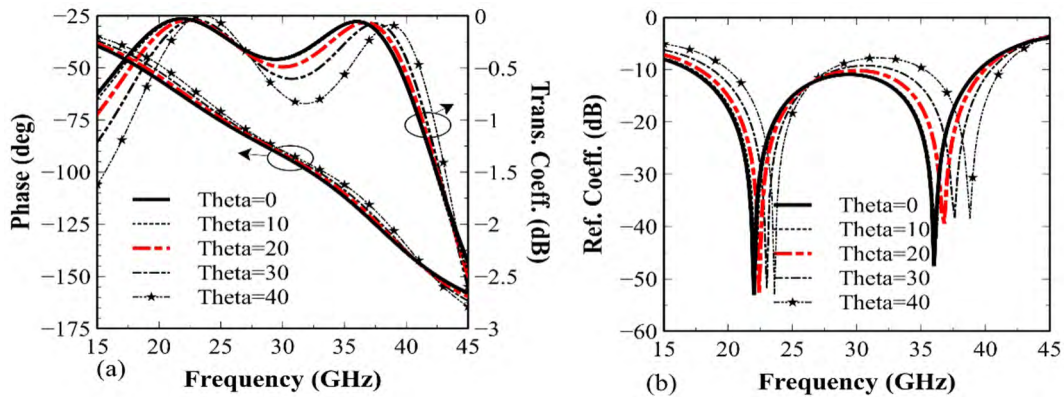
##### A. REFLECTION AND TRANSMISSION COEFFICIENTS

In order to investigate and verify the polarizer performance, a full-wave simulation is carried out using Ansys HFSS-15. Initially, the dielectric multilayer slabs with numerical values obtained from (13) to (16) are simulated. The magnitude and phase of transmission coefficients are presented in Fig. 4 (a) and (b), respectively. It can be noticed that the cross transmission coefficients  $Tr(xy)$  and  $Tr(yx)$  are nearly zero. Concerning the reflection coefficient curve depicted in Fig. 4 (c), there are two resonant frequencies at about

21GHz and 41GHz. Fig. 4(d) shows the magnitude and phase differences between transmission coefficients for both polarizations (xx and yy). It is observed that the phase difference ( $90 \pm 20^\circ$ ) and magnitude difference (less than 1dB) of transmission coefficients for both polarizations remain between 24GHz to 38GHz (45%). In terms of design simplicity without any conductors, the dielectric multilayer slabs can be a good polarizer candidate. However, the  $-10$  dB impedance bandwidth is very limited. In order to satisfy the standard impedance bandwidth ( $S_{11} \leq -10$ dB), a circular conductor, due to its shape simplicity, is printed on one side of



**FIGURE 6.** The simulated of transmission coefficients of dielectric slab with the conductor in unit cell model (a) magnitude of transmission coefficients, (b) phase of transmission coefficients, (c) reflection coefficient, and (d) difference of phase and magnitude between transmission coefficients.



**FIGURE 7.** The results related to the dielectric slab with circle conductor in the unit cell model when is illuminated by an incident wave with different oblique angles (a) magnitude and phase of transmission coefficients, (b) reflection coefficients.

each dielectric slab. The circular conductor gives the designer freedom to achieve the desirable impedance bandwidth. The dependence of transmission and reflection coefficients of the dielectric slab on the radius of the integrated circular conductor are given in Fig. 5. A conducting circle with radius 0.3mm is chosen as a reasonable value.

Finally, the simulated transmission and reflection coefficients of the proposed polarizer with the conductor in the unit cell model are depicted in Fig. 6. It is obvious that the dielectric slab with a circular conductor has the same behavior in transmission coefficient curves in comparison with the dielectric slab without the conducting circle. The main

improvement is observed in the bandwidth of reflection coefficient ( $S_{11} \leq -10\text{dB}$ ), where it covers the frequency band 22.5GHz to 40GHz (56%). However, the phase difference of transmission coefficients ( $90 \pm 20^\circ$ ) limits the frequency band from 23GHz to 35GHz (42%). All obtained results are based on normal incident EM wave when the LP wave is tilted away from the x-axis by  $45^\circ$ . Practically, the incident wave can be oblique or deflected due to the type of radiation pattern of antenna or installation inaccuracy. Therefore, the robustness of the polarizer under oblique incidence wave must be examined as well. As illustrated in Fig. 7, it is observed that the broad  $-10\text{dB}$  impedance bandwidth is still maintained up to

an oblique angle of  $30^\circ$ . It is also found out that by varying the incident wave angle, the transmission performance is approximately stable.

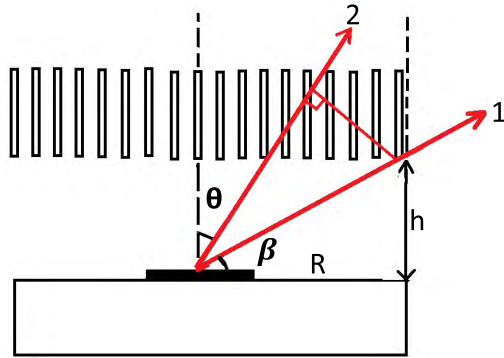


FIGURE 8. Model of the diffraction and transmission rays [2].

### B. DIELECTRIC MULTILAYER SLABS ARRANGEMENT

One of the challenges in designing multilayer structures as superstrate (placed above the radiating sources) is to specify the number of layers and estimate the size of the superstrate. To address this challenge, the model of diffraction and transmission ray is shown in Fig. 8 [2]. This technique is useful for circular layers, but it can also give a good estimation for square shape layers. In this model, the phase of diffraction and transmission rays radiated from the antenna source to dielectric multilayer slabs should be equal.

In Fig. 8, rays 1 and 2 are the first diffraction and transmission with phases 1 and 2, respectively. Since the wave phase is defined as

$$\phi = e^{-j(\beta d)} = e^{-j(\frac{2\pi}{\lambda_0} \sqrt{\epsilon_r} \cdot d)} \quad (17)$$

Therefore, with respect to Fig. 8, we have

$$(\phi)_1 = e^{-j(\frac{2\pi h}{\lambda \sin \beta})} \quad (18)$$

$$(\phi)_2 = e^{-j(\frac{2\pi h \cos(\frac{\pi}{2} - \beta - \theta)}{\lambda \sin \beta} + (\phi_t))} \quad (19)$$

In order to keep the diffraction and transmission rays in phase, we impose the following conditions,

$$\frac{2\pi h}{\lambda \sin \beta} = \frac{2\pi h \cos(\frac{\pi}{2} - \beta - \theta)}{\lambda \sin \beta} + (\phi_t) \quad (20)$$

$$\frac{h}{\lambda} \left( \frac{1}{\sin \beta} - 1 \right) - \frac{(\phi_t)}{2\pi} = N \quad (21)$$

The proposed polarizer involves the transmission phase  $\phi_t$  in x- and y- polarization, which are  $179.7^\circ$  and  $87.5^\circ$ , respectively. In order to have the minimum size of the polarizer, “N” should be a positive integer. The elevation angle  $\beta$  from the antenna source to the edge of the polarizer, for the two aforementioned directions, is computed as  $15^\circ$  and  $17^\circ$ , respectively, resulting in an ellipse with major and minor axes of 37mm and 33.4mm. Based on the diffraction and transmission rays’ model, the size of the

proposed dielectric slab as exhibited in Fig. 2 is  $33 \times 25$  unit cells. It means that the proposed polarizer consists of 33 circular conductors which are linearly printed on each of the 25 dielectric layers.

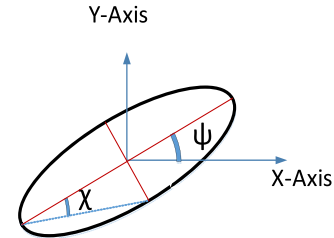


FIGURE 9. The polarization ellipse with two angles shown  $\psi$  and  $\chi$  as the geometrical parameters of the ellipse.

### C. POSITION OF THE TRANSMITTED WAVE POLARIZATION

In order to investigate the purity of transmitted wave CP, the polarization azimuth angle  $\psi$ , and ellipticity angle  $\chi$  are used to examine the orientation and case of the elliptical polarization, as shown in Fig. 9.

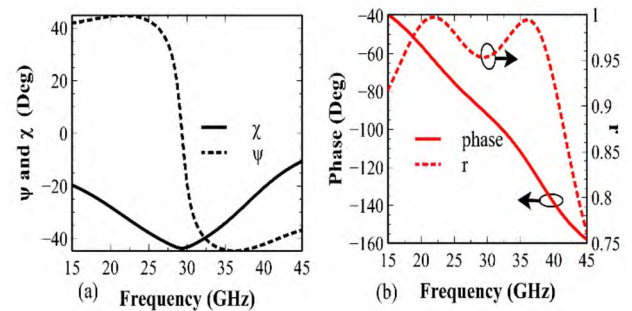


FIGURE 10. (a) The simulated angles of polarization azimuth ( $\psi$ ) and ellipticity ( $\chi$ ) and (b) the phase difference ( $\Delta\phi$ ) and amplitude ratio ( $r$ ) versus frequency.

The parameter  $\psi$  specifies the angle between the major axis of the ellipse and the x-axis, while the angle  $\chi$  examines the ellipticity behavior. The angles  $\psi$  and  $\chi$  can be calculated from the following equations [21]:

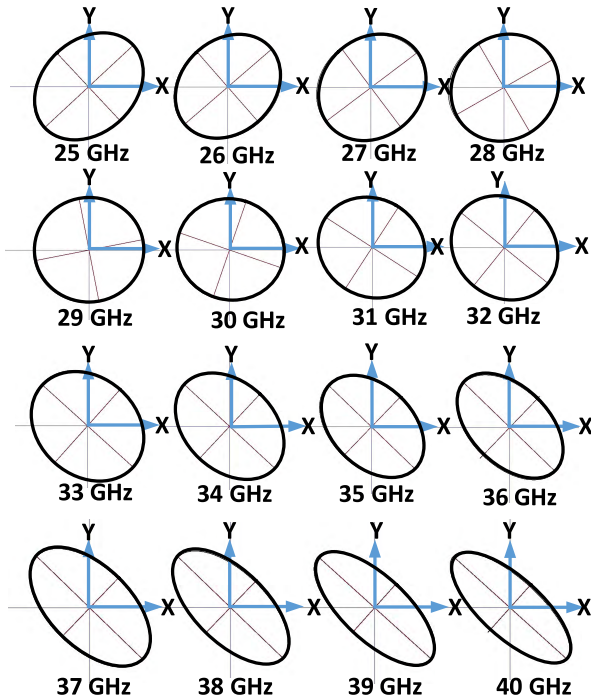
$$r = \frac{Tr(yy)}{Tr(xx)} \quad (22)$$

$$\Delta\phi = \phi_{yy} - \phi_{xx} \quad (23)$$

$$\psi = \frac{1}{2}(\tan^{-1} \left[ \frac{2r}{1-r^2} \right] \cos \Delta\phi) \quad (24)$$

$$\chi = \frac{1}{2}(\sin^{-1} \left[ \frac{2r}{1+r^2} \right] \sin \Delta\phi) \quad (25)$$

where “r” in (22) is the amplitude ratio. As shown in Fig. 10 (a), the variations of polarization azimuth angle ( $\psi$ ) and ellipticity angle ( $\chi$ ) of the transmitted wave with frequency are highlighted. Whenever the ellipticity angle of the transmitted wave is  $\pm 45^\circ$ , it means that polarization is purely circular. Correspondingly, the ellipticity angle of the transmitted wave in the proposed polarizer is  $-45^\circ$  at



**FIGURE 11.** Theoretically predicted polarization cases in the plane perpendicular to the wave vector at different frequencies from 25GHz to 40 GHz.

around 30GHz, indicating the pure CP at the frequencies close to 30GHz, as shown in Fig. 10. To further understand the tendency of the polarization ellipses and its dependence on the frequency, sixteen uniformly distributed frequencies along the bandwidth were selected. The corresponding results are shown in Fig. 11. It can be noticed that by increasing the frequency from 25GHz to 40 GHz, the polarization is right hand rotated. Also, as an acceptable estimation, it is observed that when the ellipticity angle ( $\chi$ ) corresponds to  $-45 \pm 5^\circ$ , an excellent CP from 26GHz to 33GHz is obtained. However, an elliptical polarization is observed on the rest of the frequency band. It is expected that the distortions in the circular polarization at the edges of the band are mainly caused by elevated reflections.

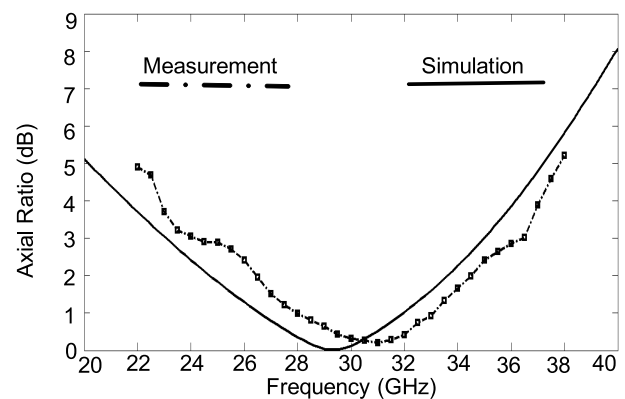
Table 1 summarized a brief comparison between the proposed design and sample published polarizers in literature in terms of 3-dB AR bandwidth and maximum insertion loss in the band of interest. According to the table, the proposed design has a better performance compared to other structures, particularly in the insertion loss performance. In order to realize the polarizer performance in practical applications, it is implemented and axial ratio is measured using a wideband horn antenna, as shown in Fig. 12. Fortunately, the measured axial ratio approximately has an acceptable agreement with the simulated one, covering 41% ( $AR \leq 3\text{dB}$ ).

## V. INTEGRATED ANTENNA DESIGN

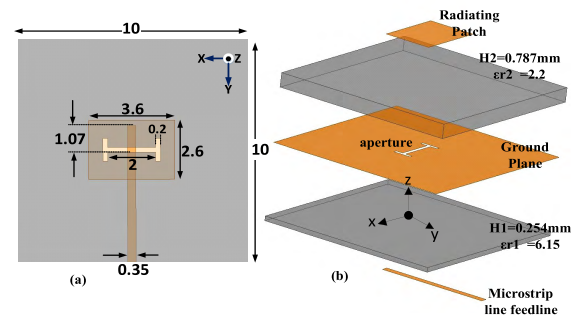
In order to investigate and verify the polarizer performance with an antenna, an 8-element linear array antenna is designed. The wide bandwidth and broadside pattern are two

**TABLE 1.** Performance comparison of our proposed design with designs in the literature.

Design	3dB AR bandwidth (%)	Peak Insertion loss (dB)
Multilayer based on two-dimensionally anisotropic FSS [5]	40	$\leq 3$
Meander line and stub [8]	36	1
Cross slot [22]	14.5	3.2
Split circular ring FSS [22]	21	3.3
This work	42	$\leq 0.5$



**FIGURE 12.** The measured and simulated axial ratio of the proposed polarizer with the wideband horn antenna.



**FIGURE 13.** The Geometry of the single linear ACMA (a) top view and (b) side view (optimized values are versus millimeter).

main reasons for choosing an aperture-coupled microstrip antenna (ACMA) as the primary prototype. Fig. 13 shows the geometry of a single ACMA feeding operating on Ku-band around 30GHz. The antenna consists of two different substrates, a Rogers RT/Duroid 5880 with dielectric constant 2.2 and thickness 0.787mm on the top and the other is Rogers 3006 with dielectric constant 6.15 and thickness 0.254mm on the bottom. Due to the limitation in the inter-element spacing of linear array structures, an H-shaped slot is etched off the common ground plate [23]. Fig. 14 illustrates an equivalent circuit of ACMA similar to the model in [24]. Fig. 15, shows the reflection coefficient and total gain of the

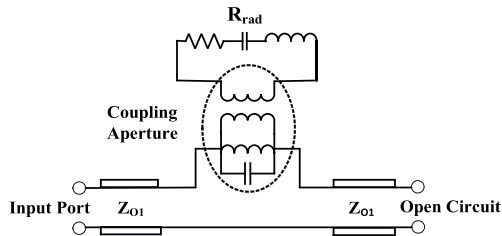


FIGURE 14. The equivalent circuit model of the ACMA [24].

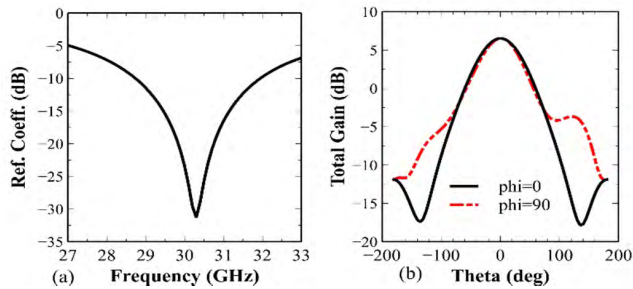


FIGURE 15. (a) The reflection coefficient and (b) radiation pattern of single ACMA at 30GHz.

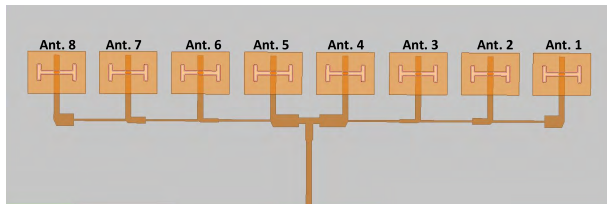


FIGURE 16. The 8-element linearly polarized array ACMA with Chebyshev tapered distribution.

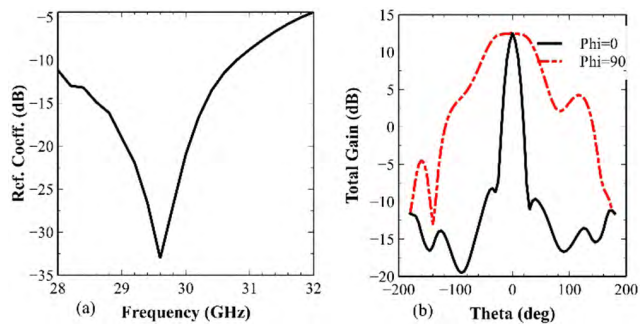


FIGURE 17. (a) The reflection coefficient and (b) radiation pattern of 8-element linear array antennas at 30GHz.

single ACMA. In order to illuminate the transmission-type polarizer, the high gain array antenna with the capability of lower sidelobe level (SLL) is proposed. Thus, an 8-element series-fed linearly polarized array antenna using suitable Chebyshev taper distribution operating at 30GHz is designed. Each antenna element is fed with a 50-ohm input impedance. To design the microstrip series-fed array, firstly the feeding network is divided into two linear subarrays and fed at the center by microstrip line [25].

In order to guarantee the same phase among the elements and to obtain the radiation pattern at broadside, the

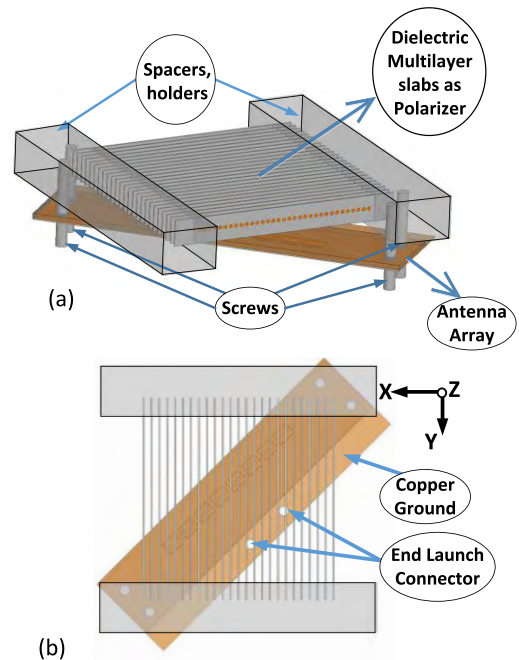


FIGURE 18. The final proposed design, including the multilayer slab polarizer, the antenna array, spacer, screws; (a) side view and (b) top view.

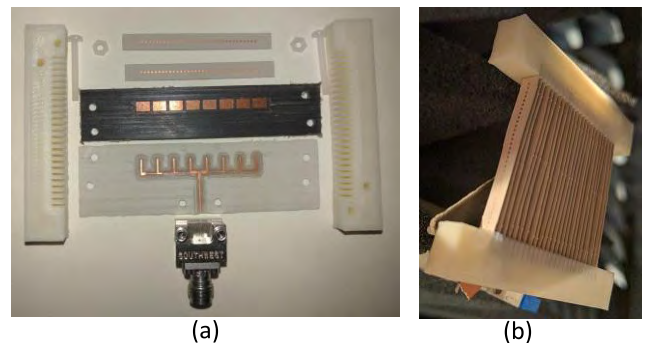
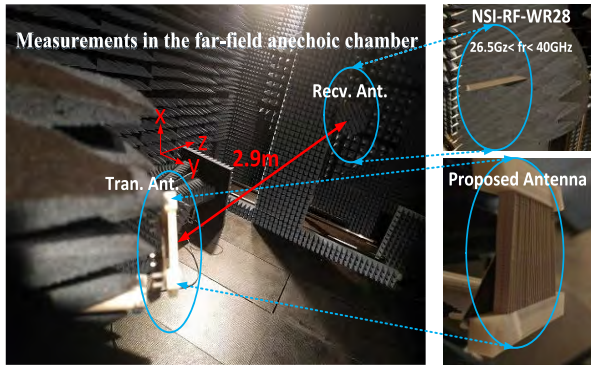


FIGURE 19. The photo of (a) the assembled design and (b) the propose structure.

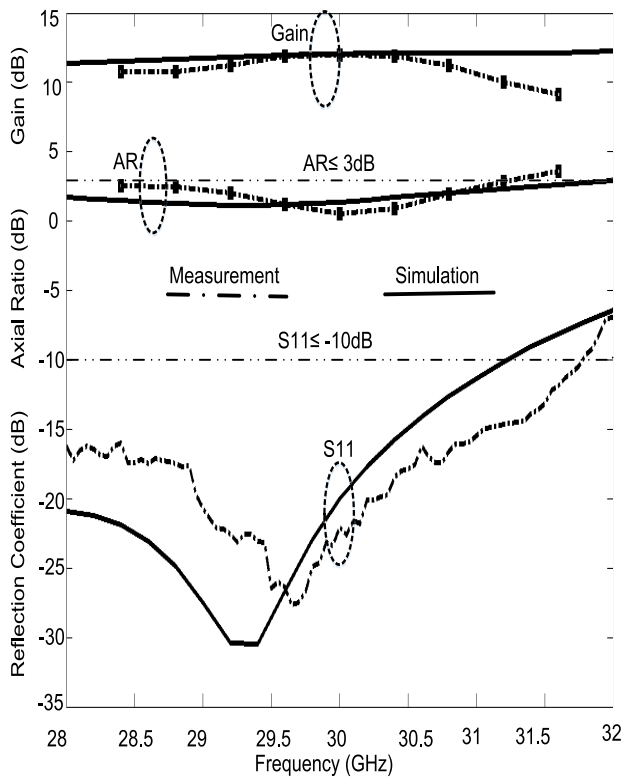
distance between the feed points of the array elements has to be equal to one guided wavelength ( $\lambda_g$ ). The symmetric arrangement reduces the cross-polarization level of the array and minimizes the beam-squinting with frequency [25], [26]. For this case, the cross-polar component produced in one direction of the antenna array is abandoned by the cross-polar component created in the reverse orientation of the antenna array at broadside. Moreover, a tapered distribution is achieved utilizing quarter-wavelength transformers ( $\lambda_g/4$ ) along the line. The simulated reflection coefficient along with radiation pattern on the E-plane and H-plane at 30GHz are shown in Fig. 17. It is observed that the ACMA array resonates at a center frequency close to 30GHz with  $-23$ dB H-plane SLL and a gain of 13 dB.

## VI. FABRICATION AND EXPERIMENTAL RESULTS

In this section, both the 8-element ACMA LP array and polarizer structures are combined to form the final proposed structure. The LP array is placed under the polarizer using

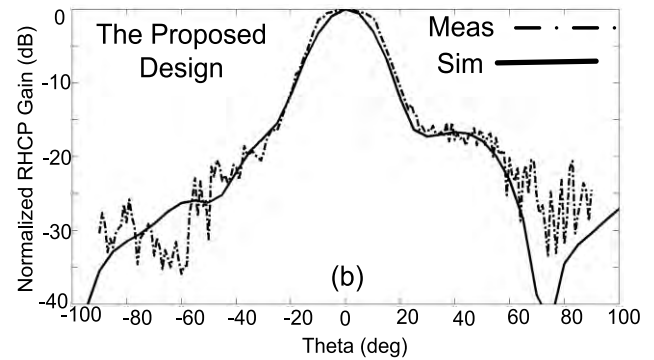
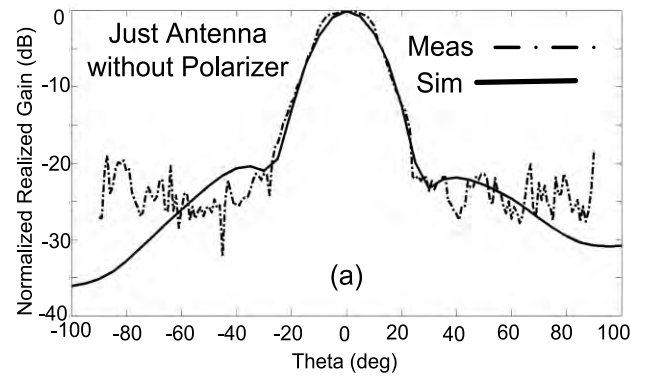


**FIGURE 20.** The photo of the antenna under test in the far-field anechoic chamber.



**FIGURE 21.** The measured and simulated curves of realized gain, axial ratio, and reflection coefficient of the proposed design; including the LP antenna array with the polarizer.

two spacers (holders) with a thickness 5mm (a half wavelength at 30 GHz), as shown in Fig. 18. In order to connect the spacers to the antenna array structure, four screws are used. In addition, an end launch connector from Southwest Company to feed and excite the array is applied. To verify the design and integrated structure's performance, the proposed dielectric multilayer slab polarizer and the 8-element LP-ACMA were fabricated and tested. Fig. 19 shows the photo of the assembled prototype. It consists of two different substrates as the antenna array, dielectric multilayer slabs as the polarizer, spacers, screws, and connector. The measured



**FIGURE 22.** (a) The normalized realized gain of the LP antenna array without the polarizer and (b) the normalized RHCP gain of the polarizer with the antenna array at 30GHz.

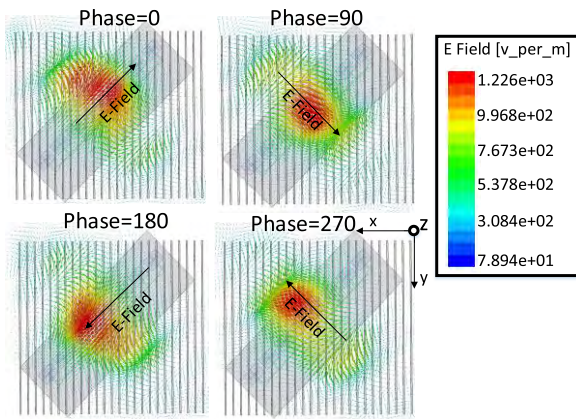
reflection coefficient of the 8-element linear array ACMA with and without the polarizer were carried out utilizing an Agilent N5227A PNA Network Analyzer (10MHz-67GHz). Also, the far field anechoic chamber is utilized to measure the radiation patterns and axial ratio as depicted in Fig. 20. The left and right-hand electric fields, and AR of the antenna are calculated from the measured electric fields, with both amplitude and phase following [27], [28]:

$$\vec{E}_{RH} = \frac{1}{\sqrt{2}} (\vec{E}_\theta + j\vec{E}_\phi) \quad (26)$$

$$\vec{E}_{LH} = \frac{1}{\sqrt{2}} (\vec{E}_\theta - j\vec{E}_\phi) \quad (27)$$

$$AR(dB) = 10 \log \left( \frac{|\vec{E}_{RH}| + |\vec{E}_{LH}|}{|\vec{E}_{RH}| - |\vec{E}_{LH}|} \right) \quad (28)$$

The measured and simulated results including the reflection coefficient, axial ratio, and realized gain are exhibited in Fig. 21. It can be observed from Fig. 21 that the experimental results and simulated ones have an acceptable agreement. The measured results of the polarizer along with the antenna array show that impedance bandwidth ( $|S_{11}| \leq -10dB$ ) covers 28GHz to 32GHz with an axial ratio bandwidth ( $AR \leq 3dB$ ) from 28GHz to 31.5GHz, and a realized gain of 13dBi with a pure CP at the central frequency of 30GHz. As shown in Fig. 22, the radiation patterns of the proposed antenna array with and without the polarizer at 30GHz are presented.



**FIGURE 23.** (a) RHCP for E-field variations at different phases (a)  $0^\circ$ , (b)  $90^\circ$ , (c)  $180^\circ$ , and (d)  $270^\circ$  at the centre frequency 30 GHz.

From the normalized realized gain curve in Fig. 22(a), it can be observed that the sidelobe level (SLL) is nearly  $-20$  dB due to the Chebyshev tapered distribution used in the feed line. Fig. 22 (b) shows the normalized RHCP gain of the final proposed integrated structure including both the antenna array and polarizer.

The simulated and measured results have an acceptable agreement in the main beam angular region. In Fig. 23, it can be realized that the radiated field from the antenna rotates in clockwise direction demonstrating a right-hand circularly polarized wave.

## VII. CONCLUSION

In this paper, a new polarizer based on dielectric multilayer slabs for MMW applications is presented. The polarizer can function from 23 GHz to 35 GHz (42%) with excellent impedance and AR bandwidths. The insertion loss of polarizer is less than 0.5 dB which in industrial applications can be excellent. The polarizer is fabricated and measured by a wideband horn antenna satisfying the simulated results. Next, in order to design a high-gain circularly polarized structure at around 30 GHz, an 8-element LP array antenna with Chebyshev tapered distribution is designed and integrated with the polarizer. Obviously, the antenna restricts the overall bandwidth due to the narrowband nature of the LP antenna array. The combined structure was fabricated and measured. There is an excellent agreement between the experimental results with the simulated ones. The measured results of the fabricated structure show that the final design is able to cover impedance and axial-ratio bandwidths nearly 28 GHz to 31.5 GHz. Furthermore, it has a pure CP and realized a gain of 13 dBi at 30 GHz. This polarizer with the antenna array can be a suitable candidate for several applications including satellite communications, orthogonal polarization transformers, sub-reflectors, and CP lenses.

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