# Model for triple-point trajectory of shock wave reflection over cylindrical 

 concave wedgeXueqiang Yuan and Jin Zhou*<br>National University of Defense Technology, Changsha, 410073, China<br>Xiaocheng Mi<br>McGill University, Montréal, Québec H3A 2K6, Canada<br>and<br>Hoi Dick Ng<br>Concordia University, Montreal, H3G 1M8, Canada

## Nomenclature

$H_{\mathrm{m}}=$ Mach stem height
$H_{\mathrm{m}}^{\text {max }}=$ maximum Mach stem height
$K_{i}=$ intersection point for triple-point trajectories of insert shock wave and Mach stem at the $i$-th wedge
$M=$ Mach number of incident shock wave
$n=$ quantity of small wedges
$O_{i}=$ tip of the $i$-th wedge
$R_{0}=$ dimensionless wall curvature radius
$T_{i}=$ Mach reflection triple-point of insert shock wave at the $i$-th wedge
$T_{\mathrm{M}}=$ Mach reflection triple-point of Mach stem

[^0]$\gamma=$ specific heat ratio
$\Delta \theta_{\mathrm{w}}=$ wedge angle difference
$\theta_{\mathrm{t}}=$ wall tangential angle
$\theta_{\mathrm{t}}^{\text {max }}=$ wall tangential angle for maximum Mach stem height
$\theta_{\mathrm{t}}^{*}=$ wall tangential angle for Mach reflection-regular reflection transition
$\theta_{\mathrm{w}}=$ planar wedge angle
$\theta_{\mathrm{w}}^{i}=$ the $i$-th wedge angle
$\theta_{\mathrm{w}-\text { shock }}^{*}=$ critical wedge angle of Mach reflection calculated by the two-shock theory
$\theta_{\mathrm{w}-\text { shock }}^{*}=$ critical wedge angle of Mach reflection calculated by the three-shock theory
$\chi=$ triple-point trajectory angle
$\chi_{i}=$ triple-point trajectory angle for insert shock wave at the $i$-th wedge
$\chi_{\mathrm{M}}=$ triple-point trajectory angle for Mach stem
$\chi_{0}=$ corresponding trajectory angle when $\theta_{\mathrm{w}}$ approaches to zero in $\chi-\theta_{\mathrm{w}}$ relation
$\chi_{2 \text { shock }}=$ corresponding trajectory angle of $\theta_{\mathrm{w}-\text {-shock }}^{*}$ in $\chi-\theta_{\mathrm{w}}$ relation

## I. Introduction

The study of unsteady shock wave reflection over curved surfaces draws much attention due to its practical importance for various applications, such as propulsion system design and aviation safety [1-4]. As the foundation of this study, the reflection problem with the simplest cylindrical concave walls has been investigated for several decades. The relevant studies were also applied for explosion and detonation ignition in recent years,
since this phenomenon involves drastic changes in pressure and temperature [5-6]. In previous research, the reflection process over a cylindrical concave wedge was described as a series of continuous variations of reflection types: direct-Mach reflection to stationary- Mach reflection to inverse-Mach reflection and finally to transitioned regular reflection [7]. To predict the trajectory of the Mach reflection triple-point and the transition angle of Mach-to-regular reflection, both the Chester-Chisnell-Whitham theory [8] and the "corner-signal" concept [9] were employed to construct models [10, 11]. Although the results indicated that the theoretical transition angle agrees with the experimental observation, the triple-point trajectory, however, cannot be predicted precisely, especially for strong shock waves with high Mach numbers.

In contrast to the shock wave reflection over a curved surface, the pseudo-steady shock reflection over a planar wedge is self-similar, so the classic two-shock and three-shock theories [12] can be adopted to predict the triple-point trajectory as well as the critical wedge angle for Mach reflection with reasonable accuracy [13]. Furthermore, the shock wave reflection over a double wedge with different conditions can also be analyzed by these theories [14]. Based on these theoretical studies, a model was constructed to predict the triple-point trajectory of strong shock wave reflection over a cylindrical concave wedge in this work. The detailed calculation process as well as a validation of the model was presented, and the characteristics of the triple-point trajectory were then analyzed by applying the model.

## II. Model setup

The two-shock and three-shock theories were first proposed by von Neumann to describe the regular reflection and the flow field near the Mach reflection triple-point, respectively [12]. By giving two constraining conditions, the three-shock theory can be used to describe the entire Mach reflection process [13]: (1) the Mach stem is straight and normal to the wedge; (2) the triple-point trajectory originates from the tip of the wedge. Since the pseudo-steady Mach reflection of a shock wave can be regarded as a steady Mach reflection by
transforming the coordinates, as Fig. 1 shows, the flow field parameters as well as the trajectory angle $\chi$ can be obtained by applying the three-shock theory with given shock wave Mach number $M$ and wedge angle $\theta_{\mathrm{w}}$. Figure 2 plots the calculated $\chi-\theta_{\mathrm{w}}$ relation for $M=5.4$ and $\gamma=1.44$, the detailed calculation process is introduced in [7]. It is worth noting that the part behind the critical wedge angle $\theta_{\mathrm{w}-3 \text { shock }}^{*}$ of Mach reflection is meaningless for pseudo-steady Mach reflection because the regular reflection is formed with the wedge angle, but the negative trajectory angle $\chi$ exists in the unsteady reflection process over a cylindrical concave wedge, which corresponds to the inverse-Mach reflection. To verify the accuracy of the theoretical $\chi-\theta_{\mathrm{w}}$ relation, the shock wave reflection process with $\theta_{\mathrm{w}}=15^{\circ}$ was simulated under the condition of $M=5.4$ and $\gamma=1.44$, and the density schlieren overlays with a red line showing the triple-point trajectory are displayed in Fig. 3. The theoretical trajectory angle $\chi$ for $\theta_{\mathrm{w}}=15^{\circ}$ is $14.12^{\circ}$, approaching to the numerical result of $15.21^{\circ}$.


Figure 1. Schematic of pseudo-steady Mach reflection structure over a planar wedge.


Figure 2. $\chi-\theta_{\mathrm{w}}$ relation for $M=5.4$ and $\gamma=1.44$.
( $\theta_{\mathrm{w}-3 \text { shock }}^{*}$ : Critical wedge angle of Mach reflection by the three-shock theory)


Figure 3. Density schlieren overlays of shock wave reflection over a planar wedge with $\theta_{\mathrm{w}}=15^{\circ}$.
Based on the pseudo-steady reflection and $\chi-\theta_{\mathrm{w}}$ relation, the reflection structure and the triple-point trajectory over a double planar wedge can also be predicted [14]. Figure 4 displays the reflection process over a double wedge under the condition of $\theta_{\mathrm{w}}^{2}>\theta_{\mathrm{w}}^{1}$. As Fig. 3(a) shows, when the incident shock wave propagates to wedge $A$ with wedge angle $\theta_{\mathrm{w}}^{1}$, the Mach reflection is formed and triple-point $T_{1}$ moves along the straight trajectory with angle $\chi_{1}$. As the Mach stem reaches wedge $B$ along with the incident shock, another Mach reflection is also formed on the Mach stem, generating the triple-point $T_{\mathrm{M}}$. Since the Mach stem is perpendicular to wedge $A$, the newly formed triple-point trajectory angle $\chi_{\mathrm{M}}$ actually corresponds to the wedge angle difference $\Delta \theta_{\mathrm{w}}$. Then the trajectories of triple-points $T_{1}$ and $T_{\mathrm{M}}$ finally intersect at point $K_{1}$, as shown in Fig. 4(b). A single Mach stem perpendicular to wedge B is formed, and the triple-point will then propagate along the trajectory with the angle of $\chi_{2}$ generated by the wedge with the angle of $\theta_{\mathrm{w}}^{2}$. According to the reflection process, the triple-point trajectory over the double wedge can be predicted by obtaining the trajectory angles $\chi_{1}, \chi_{\mathrm{M}}$ and $\chi_{2}$ and determining the location of the intersection point $K_{1}$.

(a) Mach reflection on incident shock wave and Mach

(b) Mach reflections intersect
stem simutaneously

Figure 4. Schematic of reflection process over a double planar wedge with $\theta_{\mathrm{w}}^{2}>\theta_{\mathrm{w}}^{1}$.

For a shock wave reflection over a cylindrical concave wedge, the wedge can be approximated by a larger number of small planar wedges with equivalent length, as displayed in Fig. 5, so the original reflection process is transformed to the reflection over a multiple segments wedge. From the reflection over a double wedge, it is inferred that the triple-point trajectory over a multiple segments wedge can be obtained by determining the location of $K$ for each adjacent double wedge. For instance, by defining that the tip of the $i$-th small wedge has the coordinates of $O_{i}\left(x_{i}, y_{i}\right)$, the coordinates of the first intersection $K_{1}\left(x_{k 1}, y_{k 1}\right)$ can be obtained by the simultaneous equations of the incident shock and Mach stem triple-point trajectories:

$$
\left\{\begin{array}{l}
y-y_{1}=\chi_{1}\left(x-x_{1}\right)  \tag{1}\\
y-y_{2}=\chi_{\mathrm{M}}\left(x-x_{2}\right)
\end{array}\right.
$$

Then the $i$-th intersection $K_{i}\left(x_{k i}, y_{k i}\right)$ can be confirmed by:

$$
\left\{\begin{array}{l}
y-y_{k(i-1)}=\chi_{i}\left(x-x_{k(i-1)}\right)  \tag{2}\\
y-y_{i+1}=\chi_{\mathrm{M}}\left(x-x_{i+1}\right)
\end{array}\right.
$$

where $\chi_{i}$ represents to the triple-point trajectory angle for the insert shock wave at the $i$-th wedge, which corresponds to the $i$-th wedge angle $\theta_{\mathrm{w}}^{i}$ in the $\chi-\theta_{\mathrm{w}}$ relation. Hence, if the quantity of the constituent small wedges is sufficient, the triple-point trajectory over a cylindrical concave wedge can be calculated approximately by confirming all the coordinates of $K_{i}$. It is worth noting that, since $\chi_{M}$ is the trajectory angle
of the Mach stem, theoretically the calculation of $\chi_{M}$ should use the Mach number of the Mach stem instead of the incident shock. However, considering that the deviation of the trajectory angle formed by the Mach stem and the incident shock is minimal under such a high Mach number, the $\chi_{\mathrm{M}}$ can be approximated by using $M$ in the calculation. Furthermore, since the concave wedge is divided equally by the small incremental wedges, the $\Delta \theta_{\mathrm{w}}$ between each two adjacent wedges is equal, resulting a constant $\chi_{\mathrm{M}}$ for all the composed double wedges, thus simplifying the calculation.


Figure 5. Schematic of cylindrical concave wedge approximated by multiple wedges with equivalent length.

## III. Results and discussions

## A. Theoretical triple-point trajectory and validation

A model to predict the trajectory of the reflection triple-point can be built by MATLAB based on the analysis above. To verify the accuracy of the model, a simulation of shock wave reflection over a cylindrical wedge was conducted with $M=5.4, \gamma=1.44$ and wall curvature radius $R_{0}=750$. The density schlieren of the reflection process is shown in Fig. 6. The trajectory was recorded by tracking the coordinates of the triplepoint with constant time intervals. Simultaneously, the theoretical trajectory was calculated by the model under the same condition with the quantity of small wedges $n=2000$. The comparison of the trajectories as well as the converted Mach stem height $H_{\mathrm{m}}$ are presented in Fig. 7. It can be observed that the trajectories converge
in the range of the wall tangential angle $\theta_{\mathrm{t}}<52^{\circ}$, whereas the $H_{\mathrm{m}}$ calculated by the model deviates from the simulation result with larger value for increasing $\theta_{\mathrm{t}}$. Additionally, the calculated trajectory does not realize Mach-to-Regular (MR-RR) transition, which is unreasonable. The reason is supposed that the $\chi-\theta_{\mathrm{w}}$ relation inferred by the three-shock theory is imprecise for large $\theta_{\text {w }}$. Indeed, the previous research has reported that the calculated $\chi$ becomes inaccurate compared with the experiment results as $\theta_{\mathrm{w}}$ increases [15], and the study of Hornung claimed that the critical wedge angle of the Mach reflection, i.e., the corresponding wedge angle of $\chi=0$, approaches to the result calculated by two-shock theory instead of three-shock theory, where the critical wedge angle $\theta_{\mathrm{w}-2 \text { shock }}^{*}$ and $\theta_{\mathrm{w}-\text {-shock }}^{*}$ for these two theories can be obtained by plotting the shock-polar curves of the incident and reflected shock [16]. Simulations were also conducted to verify the report of Hornung. The shock wave reflections with $\theta_{\mathrm{w}}=50^{\circ}$ and $52^{\circ}$ were simulated, and the results are shown in Fig. 8. It is found that a Mach reflection is still formed in the case of $\theta_{\mathrm{w}}=50^{\circ}$, while as $\theta_{\mathrm{w}}$ increases to $52^{\circ}$, a regular reflection is formed, indicating that the critical wedge angle is between $50^{\circ}$ to $52^{\circ}$. Since the $\theta_{\mathrm{w}-2 \text { shock }}^{*}$ and $\theta_{\mathrm{w}-3 \text { shock }}^{*}$ for $M=5.4$ and $\gamma=1.44$ are $51.14^{\circ}$ and $65.51^{\circ}$, respectively, it is confirmed that the critical wedge angle is close to the theoretical $\theta_{\mathrm{w}-2 \text { shock }}^{*}$ instead of $\theta_{\mathrm{w}-3 \text { shock }}^{*}$, which agrees with Hornung's study.


Figure 6 . Density schlieren overlays of shock wave reflection over a cylindrical concave wedge.


Figure 7. Comparison of triple-point trajectory and $H_{\mathrm{m}}$ for simulation and model.


Figure 8. Shock wave reflection over planar wedges with $\theta_{\mathrm{w}}$ close to the critical wedge angle.

According to the studies above, the $\chi-\theta_{\mathrm{w}}$ relation needs to be modified to insure the $\chi$ is accurate for small wedge angle and equal to zero when $\theta_{\mathrm{w}}=\theta_{\mathrm{w}-\text {-shock }}^{*}$. A simple modification of adding a correction term of $-\frac{\chi_{\text {shhock }}}{\theta_{\mathrm{w} \text {-shock }}} \theta_{\mathrm{w}}$ can meet the requirement, where $\chi_{\text {shock }}$ represents the corresponding trajectory angle of $\theta_{\mathrm{w}-\text { shock }}^{*}$ in the original $\chi-\theta_{\mathrm{w}}$ relation. For instance, the correction term for $M=5.4$ and $\gamma=1.44$ is calculated to be $-\frac{2.8}{51.2} \theta_{\mathrm{w}}$. To verify the modification, the simulations of shock wave reflection over planar wedges was conducted with various $\boldsymbol{\theta}_{\mathrm{w}}$ under the conditions of $M=5.4$ and $\gamma=1.44$. The modified $\chi$ - $\boldsymbol{\theta}_{\mathrm{w}}$
relation as well as the trajectory angle $\chi$ for each case are shown in Fig. 9. The numerical results are coincident with the relation curve, which indicates that the modification is valid. The trajectory and $H_{\mathrm{m}}$ calculated by the modified $\chi-\theta_{\mathrm{w}}$ relation are plotted in Fig. 10. It can be observed that the modified theoretical trajectory agrees well with the numerical result in the entire range. The calculated MR-RR transition angle is $\theta_{\mathrm{t}}^{*}=79.93^{\circ}$, which has a percentage error of $4.55 \%$ compared with the numerical result of $\theta_{\mathrm{t}}^{*}=76.45^{\circ}$, and the maximum percentage error of $H_{\mathrm{m}}$ is $6.35 \%$, indicating the modification improves the accuracy of the model significantly, which can be applied to predict the triple-point trajectory of the shock wave reflection over the cylindrical concave wedge. It is worth noting that similar with the previous models [10, 11], the present model cannot predict the effect of the wall curvature radius $R_{0}$. In fact, the effect of $R_{0}$ on the reflection process is controversial, and the mechanism is still uncertain so far.


Figure 9. Comparison of the modified $\chi-\theta_{\mathrm{w}}$ relation and the numerical results with $M=5.4$ and $\gamma=1.44$.


Figure 10. Comparison of triple-point trajectory and $H_{\mathrm{m}}$ for simulation and modified model.

## B. Characteristics of triple-point trajectory

Using the present model to predict accurately the shock reflection over a cylindrical concave wedge, the characteristics of the triple-point trajectory can be analyzed. Figure 11 shows the $H_{\mathrm{m}}-\theta_{\mathrm{t}}$ relations calculated by the model for different $M$. It can be observed that the $H_{\mathrm{m}}$ for the same $\theta_{\mathrm{t}}$ and the transition angle $\theta_{\mathrm{t}}^{*}$ decrease along with the increase of $M$, but the deviation is negligible, which coincides with the previous experimental study [11]. Thus, it appears that for the shock wave reflection with high Mach number, the triplepoint trajectories for different $M$ can actually be approximated by one trajectory with certain accuracy, which can simplify the analysis with varying $M$.


Figure 11. Theoretical $H_{\mathrm{m}}-\theta_{\mathrm{t}}$ relations for different $M$.

The model also indicates that the triple-point trajectory is not entirely a curve. In fact, the triple-point propagates along a straight line at the beginning of the reflection process. Figure 12 presents the trajectory in the range of $\theta_{\mathrm{t}}<30^{\circ}$ with $M=5.4$ and $R_{0}=750$. The data points calculated by the model represent the intersection points $K$ of the triple-point trajectories generated by the Mach reflection of the incident shock wave and the Mach stem over each double wedge. It is found that the distance between the wedge tip $O_{1}(0,0)$ and the first data point $K_{1}$ is much longer than that between each adjacent two data points. Hence, according to the reflection over a double wedge, the triple-point trajectory before $K_{1}$ is a straight line with the slope of $\tan \chi_{0}$, where $\chi_{0}$ denotes the trajectory angle when $\theta_{\mathrm{w}}$ approaches to zero in the $\chi-\theta_{\mathrm{w}}$ relation. Moreover, it is also found that the corresponding wall tangential angle of $K_{1}$ is always larger than $\chi_{0}$. Considering that the $H_{\mathrm{m}}-\theta_{\mathrm{t}}$ relation curve only has one maximum point, it can be confirmed that the location of the point with the maximum height of Mach stem $H_{\mathrm{m}}^{\max }$ is on the straight part of the trajectory, corresponding to the wall tangential angle of $\theta_{\mathrm{t}}^{\max }=\chi_{0}$. Thus the $H_{\mathrm{m}}^{\max }$ can also be obtained to be $H_{\mathrm{m}}^{\max }=R_{0}\left(1-\cos \chi_{0}\right)$. As a verification, the $\theta_{\mathrm{t}}^{\max }$ and $H_{\mathrm{m}}^{\max }$ were calculated for $M=5.4$ and $R_{0}=750$ with the results of $24.1^{\circ}$ and 63.37. The theoretical results have percentage errors of $1.86 \%$ and $1.15 \%$ compared with the numerical results of $\theta_{\mathrm{t}}^{\max }=23.66$ and $H_{\mathrm{m}}^{\max }=64.63$, respectively, which presents high accuracy. Hence, the $\theta_{\mathrm{t}}^{\max }$ and $H_{\mathrm{m}}^{\max }$
can be obtained just by the $\chi-\theta_{\mathrm{w}}$ relation instead of the model.


Figure 12. Triple-point trajectory in the range of $\theta_{\mathrm{t}}<30^{\circ}$.

## IV. Conclusions

In the present study, a theoretical model to predict the triple-point trajectory of strong shock wave Mach reflection over a cylindrical concave wedge was constructed based on the reflection over a double wedge and the three-shock theory. Together with a modification by applying the two-shock theory, the triple-point trajectory calculated by the model for the entire shock reflection is validated with the simulation result. The model denotes that the triple-point trajectory as well as the transition angle of Mach reflection to regular reflection are almost invariable with high Mach number, and the trajectory is straight at the beginning of the reflection process with corresponding trajectory angle when the wedge angle approaches to zero. The point with the maximum height of Mach stem is confirmed on the straight part of the trajectory, and the location as well as the maximum Mach stem height is predicted accurately.

## Acknowledgements

This work is supported by the National Natural Science Foundation of China under Grant Nos. 91541103 and 51776220 .

## References

[1] Geva, M., Ram, O., and Sadot, O., "The non-stationary hysteresis phenomenon in shock wave reflections," Journal of Fluid Mechanics, Vol. 732, 2013. https://doi.org/10.1017/jfm.2013.423
[2] Ram, O., Geva, M., and Sadot, O., "High spatial and temporal resolution study of shock wave reflection over a coupled convex-concave cylindrical surface," Journal of Fluid Mechanics, Vol. 768, 2015, pp. 219239. https://doi.org/10.1017/jfm.2015.80
[3] Soni, V., Hadjadj, A., Chaudhuri, A., and Ben-Dor, G., "Shock-wave reflections over double-concave cylindrical reflectors," Journal of Fluid Mechanics, Vol. 813, 2017, pp. 70-84. https://doi.org/10.1017/jfm. 2016.825
[4] Yuan, X., Zhou, J., Mi, X., and Ng, H. D., "Numerical study of cellular detonation wave reflection over a cylindrical concave wedge," Combustion and Flame, Vol. 202, 2019, pp. 179-194. https://doi.org/10.1016/j.combustflame.2019.01.018
[5] Skews, B. W., and Kleine H., "Flow features resulting from shock wave impact on a cylindrical cavity," Journal of Fluid Mechanics, Vol. 580, 2007, pp. 481-493. https://doi.org/10.1017/S0022112007005757
[6] Shadloo, M. S., Hadjadj, A., and Chaudhuri, A. "On the onset of postshock flow instabilities over concave surfaces." Physics of Fluids, Vol. 26, No. 7, 2014, pp. 076101. https://doi.org/10.1063/1.4890482
[7] Ben-Dor, G., "Shock wave reflection phenomena, 2nd ed.," Springer-Verlag, New York, US, 2007. https://xs.scihub.ltd/https://doi.org/10.1007/978-3-540-71382-1
[8] Whitham, G. B., "A new approach to problems of shock dynamics. Part 1. Two-dimensional problems," Journal of Fluid Mechanic, Vol. 2, 1957, pp. 146-171. https://doi.org/10.1017/S002211205700004X
[9] Hornung, H. G., Oertel, H., and Sandeman, R. J., "Transition to Mach reflexion of shock waves in steady
and pseudosteady flow with and without relaxation," Journal of Fluid Mechanics, Vol. 90, No. 3, 1979, pp. 541-560. https://doi.org/10.1017/S002211207900238X
[10] Itoh, S., Okazaki, N., and Itaya, M., "On the transition between regular and Mach reflection in truly nonstationary flows," Journal of Fluid Mechanics, Vol. 108, 1981, pp. 383-400. https://doi.org/10.1017/S0022112081002176
[11] Ben-Dor, G., and Takayama, K., "Analytical prediction of the transition from Mach to regular reflection over cylindrical concave wedges," Journal of Fluid Mechanics, Vol. 158, 1985, pp. 365-380. https://doi.org/10.1017/S0022112085002695
[12] von Neumann, J., "Collected works, 6, Pergamon," Oxford, UK, 1961.
[13] Law, C. K., and Glass, I. I., "Diffraction of Strong Shock Waves by a Sharp Compressive Corner," C.A.S.I Transactions, Vol. 4, No. 1, 1971, pp. 2-12.
[14] Ben-Dor, G., Dewey, J. M., and Takayama, K., "The reflection of a plane shock wave over a double wedge," Journal of Fluid Mechanics, Vol. 176, 1987, pp. 483-520. https://doi.org/10.1017/S0022112087000776
[15] Milton, B. E., "Mach reflection using ray-shock theory," AIAA Journal, Vol. 13, No. 11, 1975, pp. 15311533. https://doi.org/10.2514/3.60566
[16] Hornung, H., "Regular and Mach reflection of shock waves," Annual Review of Fluid Mechanics, Vol. 18, No. 1, 1986, pp. 33-58. https://doi.org/10.1146/annurev.fl.18.010186.000341


[^0]:    ${ }^{1}$ Doctoral Student, National University of Defense Technology, No. 109, Deya Street, Changsha, 410073, China.
    ${ }^{2}$ Professor, National University of Defense Technology, No. 109, Deya Street, Changsha, 410073, China, AIAA member.
    ${ }^{3}$ Doctoral Student,
    ${ }^{4}$ Professor,

